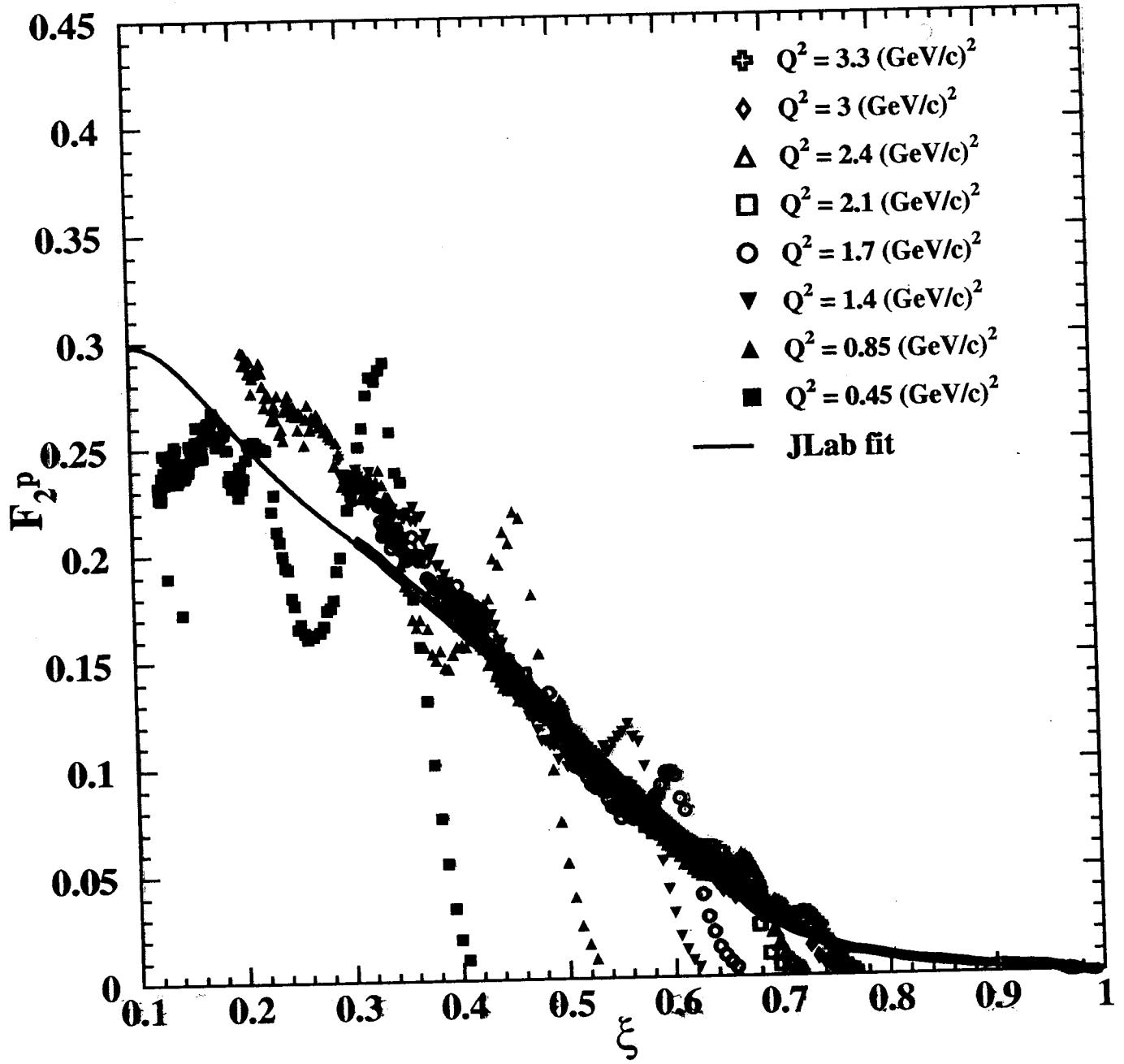
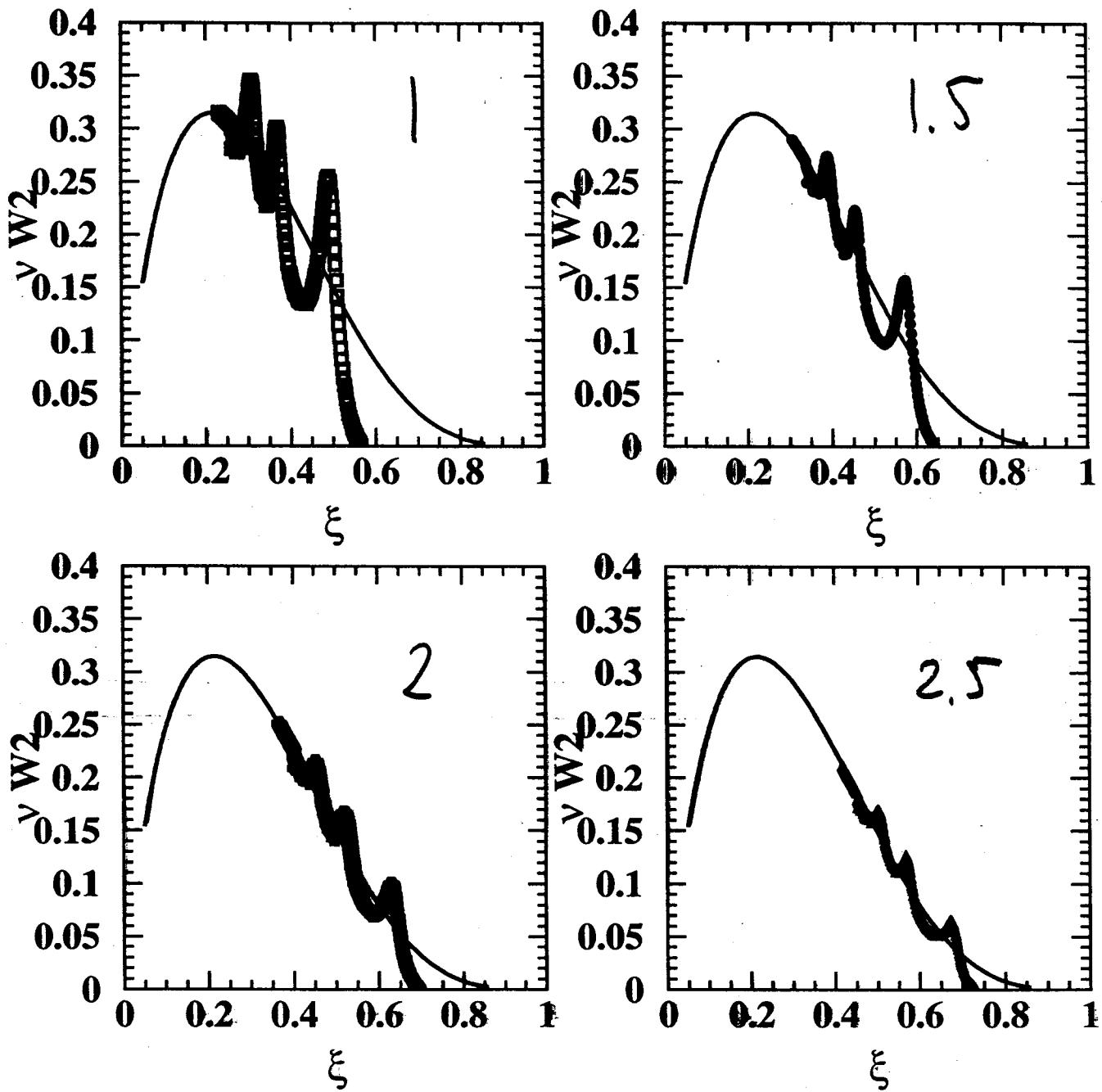


DUALITY FOR POLARIZED
AND UNPOLARIZED
STRUCTURE FUNCTIONS

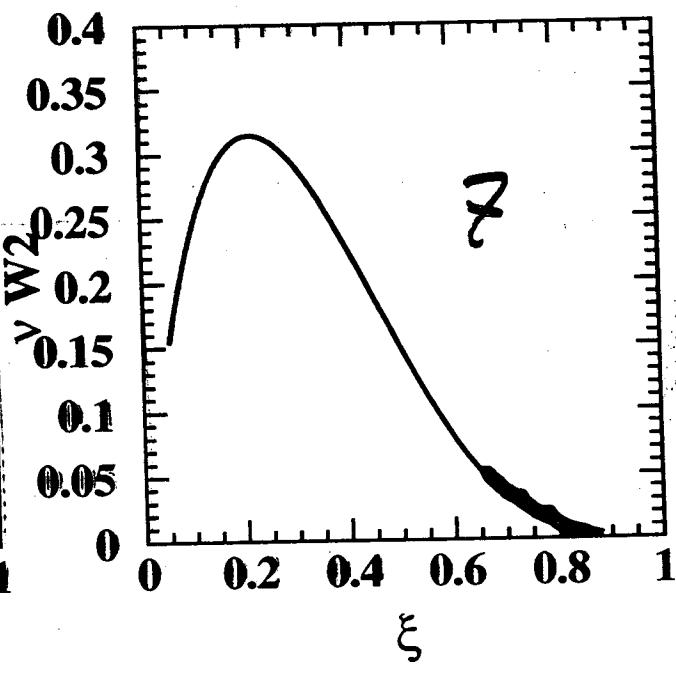
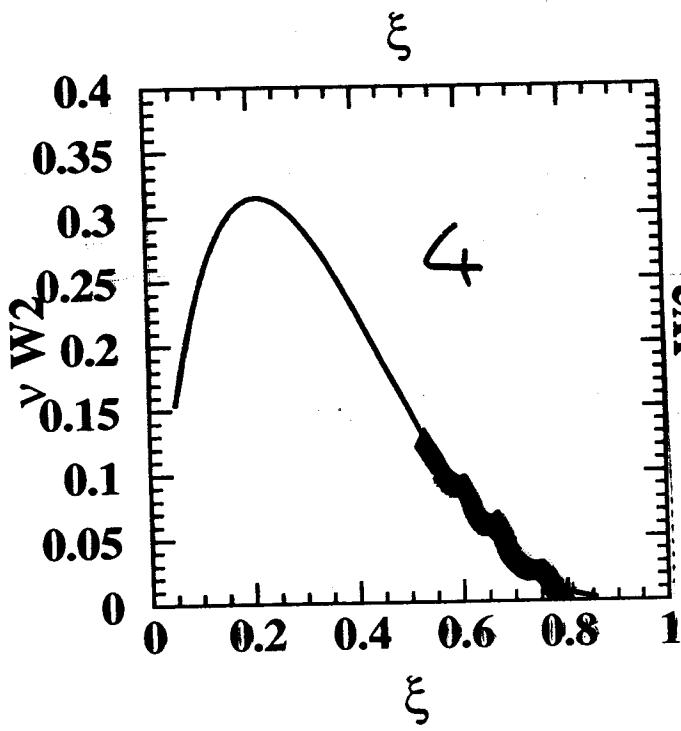
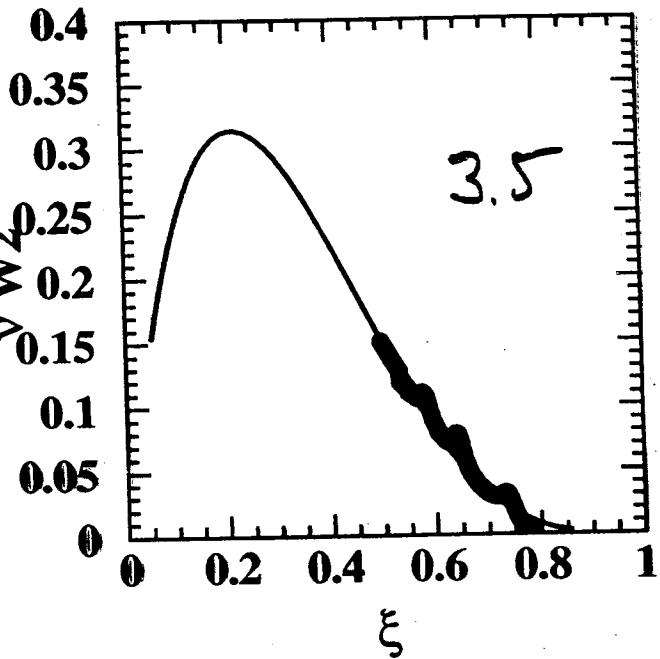
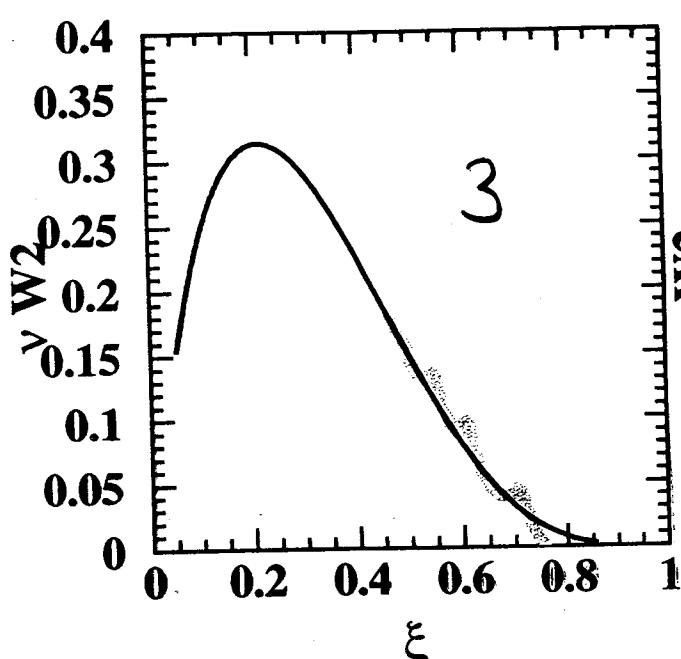
A. RADYUSHKIN



JLab Results



JLab Results

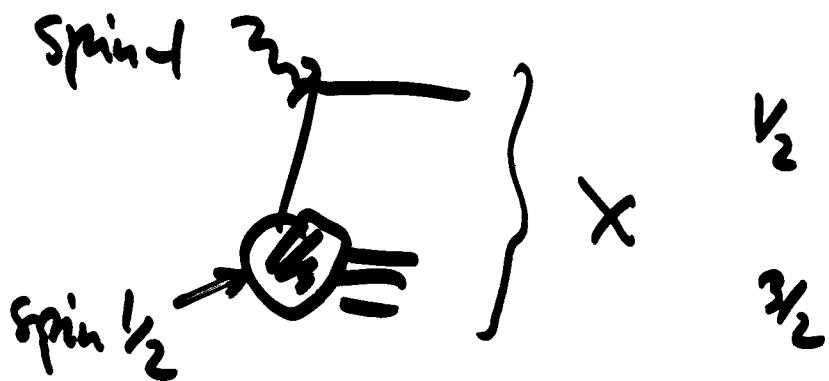


QUESTION: HOW DUALITY WILL
WORK IN POLARIZED CASE?

PROBLEM: $G_1(x, Q^2)$ IS NOT ALWAYS
POSITIVE

$$W_2 \sim \sigma^{\frac{1}{2}} + \sigma^{\frac{3}{2}}$$

$$G_1 \sim \sigma^{\frac{1}{2}} - \sigma^{\frac{3}{2}}$$



IN GENERAL: IS DUALITY POSSIBLE
IF FUNCTION IS NOT POSITIVE
DEFINITE?

CONSIDER SIMPLE CASES

DUALITY IN $e^+e^- \rightarrow \text{hadrons}$

$$\left| \sum_{\text{hadrons}} \right|^2 \sim \sigma^{\text{tot}}(e^+e^- \rightarrow \text{hadrons})$$

$$\left| \sum_q \right|^2 \sim \sum_q \left| \sum_{\text{hadrons}} \right|^2 \stackrel{\text{all possible hadrons}}{\sim} \text{Im}$$

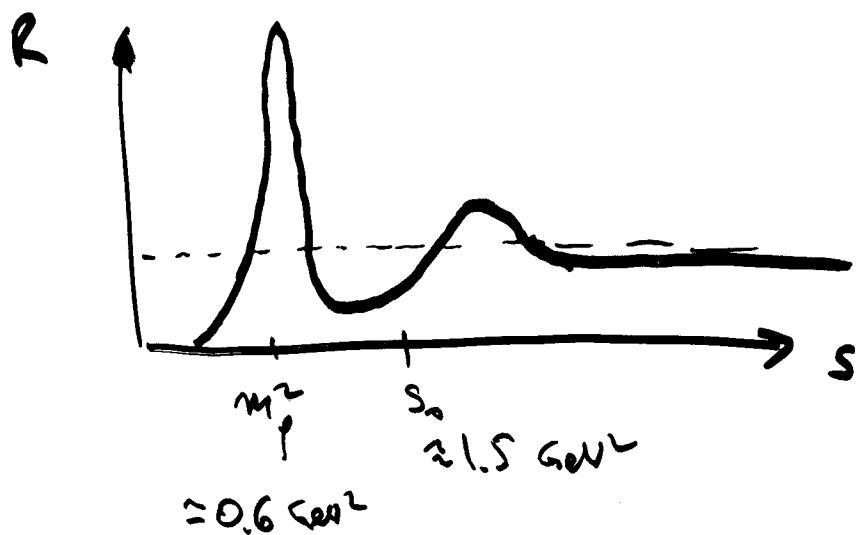
FOR LARGE $q^2 \equiv s$

$$\sum_{\substack{\dots \\ \text{HADRONS}}} = \sum_{\substack{\dots \\ \text{QUARKS} \\ \times \text{GLUONS}}}$$

$$\sum_{\substack{\dots \\ \text{HADRONS}}} \stackrel{\text{?}}{=} \sum_{\substack{\dots \\ \text{QUARKS}}} + \sum_{\substack{\dots \\ \text{GLUONS}}} + \sum_{\substack{\dots \\ \text{HADRONIC T.}}}$$

QUARKS & GLUONS WITH "100% PROBABILITY"
CONVERT INTO HADRONS (?)

$$R = \frac{\sigma^{\text{tot}}(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$



NB : FOR $s \sim m_p^2$: $\sigma^{\text{hadr}} \gg \sigma^{\text{quark}}$
 FOR SMALL s : $\sigma^{\text{hadr}} \ll \sigma^{\text{quark}}$

$$\int_0^{s_0} \sigma^{\text{hadr}}(s) ds \approx \int_0^{s_0} \sigma^{\text{quark}}(s) ds$$

\uparrow
DOMINATED BY ρ

$s_0 \approx 1.5 \text{ GeV}^2$

ρ IS DUAL TO PART OF
QUARK "SPECTRUM"

QCD SUM RULES (SVZ) FREE QUARKS

$$\int_0^\infty e^{-s/m^2} R^{I=1}(s) ds = \frac{3}{2} m^2 \left[1 + \frac{\alpha_s(m)}{\pi} - \frac{2\pi^2 f_\pi^2 m_\pi^2}{m^4} \right]$$

$$+ \frac{\pi^2}{3m^4} \underbrace{\left\langle \frac{\alpha_s}{\pi} Q\bar{Q} \right\rangle}_{\text{GLUON}} - \frac{448\pi^3}{m^6} \alpha_s \langle q\bar{q} \rangle^2 \underbrace{\text{CONDENSATE}}_{\text{QUARK}}$$

$$\approx \frac{3}{2} m^2 \left[1 + \frac{\alpha_s(m)}{\pi} + 0.1 \left(\frac{0.6 \text{ GeV}^2}{m^2} \right)^2 - 0.14 \left(\frac{0.6 \text{ GeV}^2}{m^2} \right)^4 \right] \uparrow \text{POWER CORRECTIONS}$$

(HIGHER DIMENSIONS)

TAKE $m^2 \rightarrow \infty$:

$$\int_0^\infty (R^{\text{hadr}}(s) - R^{\text{quark}}(s)) ds e^{-\frac{s}{m^2}} = \Theta\left(\frac{1}{m^4}\right) + \dots$$

$\frac{3}{2}$ 1 $\rightarrow 0$

$$\int_0^\infty (R^{\text{hadr}}(s) - R^{\text{quark}}(s)) ds = 0.$$

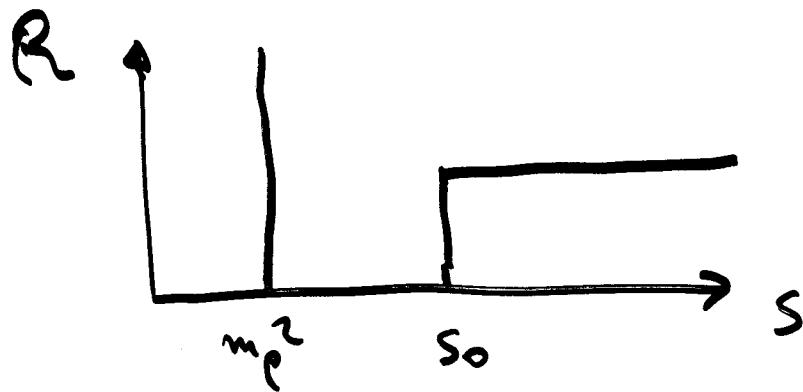
OR
GLOBAL DUALITY

$$\int_0^{s_{\text{cont}}} ds R^{\text{hadr}}(s) = \int_0^{s_{\text{cont}}} R^{\text{quark}}(s) ds$$

$s_{\text{cont}} > \text{onset} \circledast \text{of continuum}$

APPROXIMATION:

$$R^{\text{hadr}}(s) \approx \pi f_\rho^2 \delta(s - m_\rho^2) + R^{\text{quark}}(s) \Theta(s - s_0)$$



FITTING WITH KNOWN CONDENSATE

$$\text{VALUES : } m_\rho^2 \approx 0.6 \text{ GeV}^2$$

$$s_0 \approx 1.5 \text{ GeV}^2$$

SINCE $R^{\text{model}}(s) = R^{\text{quark}}(s)$ FOR $s > s_0$,

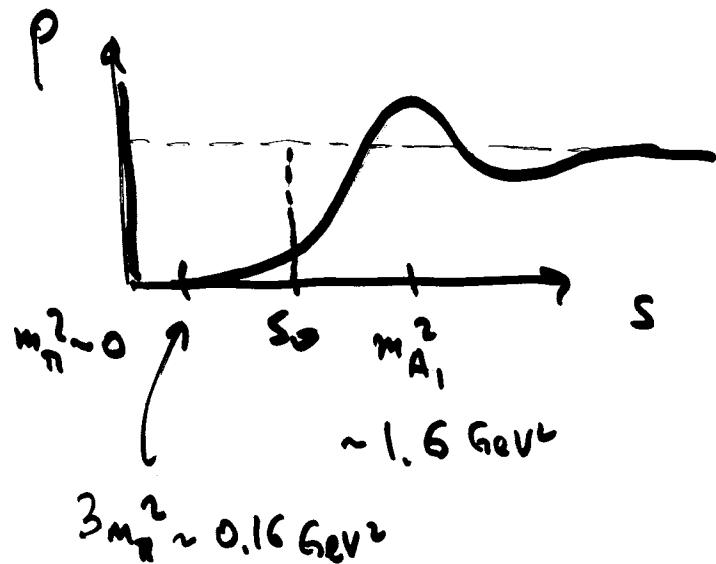
$$s_0$$

$$\int_0^{s_0} R^{\text{quark}}(s) ds = f_\rho^2$$

LOCAL DUALITY

$s_0 \equiv \rho$ duality INTERVAL

SPECTRUM IN AXIAL CURRENT CHANNEL



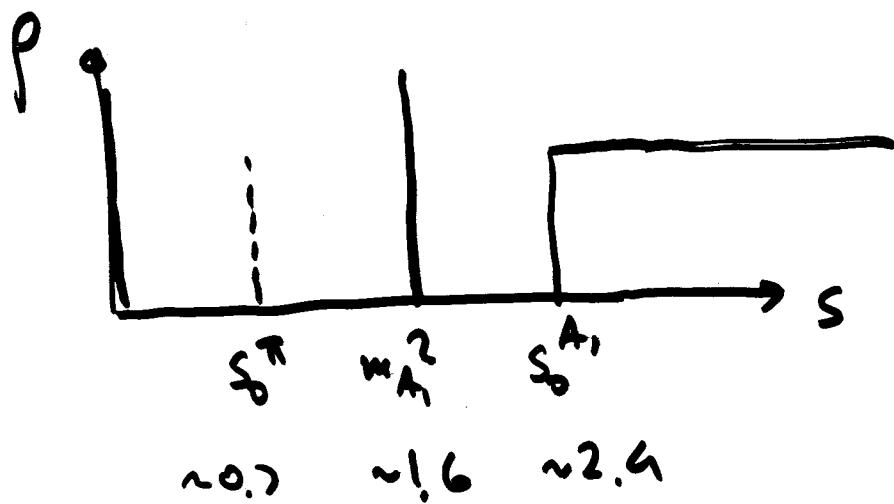
DUALITY INTERVAL:

$$s_0 = 4\pi^2 f_\pi^2 \approx 0.67 \text{ GeV}^2$$

3 π THRESHOLD

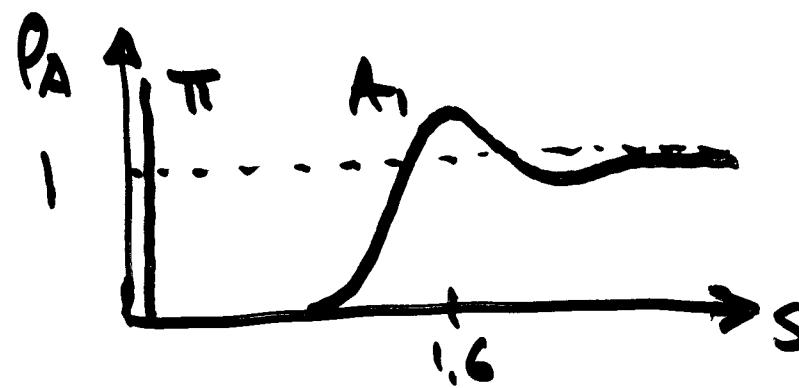
$$(3m_{\pi})^2 = 0.16 \text{ GeV}^2$$

DUALITY INTERVAL $s_0^{(\pi)}$ HAS
NOTHING TO DO WITH 3 π THRESHOLD

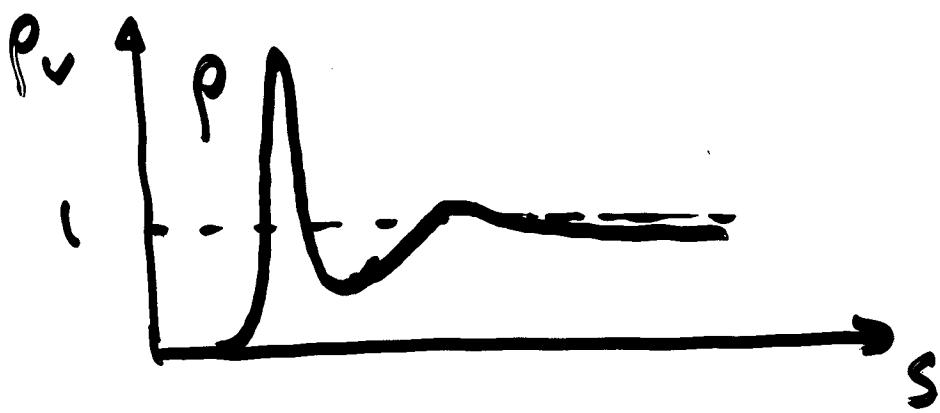


model WITH
INFINITELY
NARROW A_1 ,
(QCD SR)

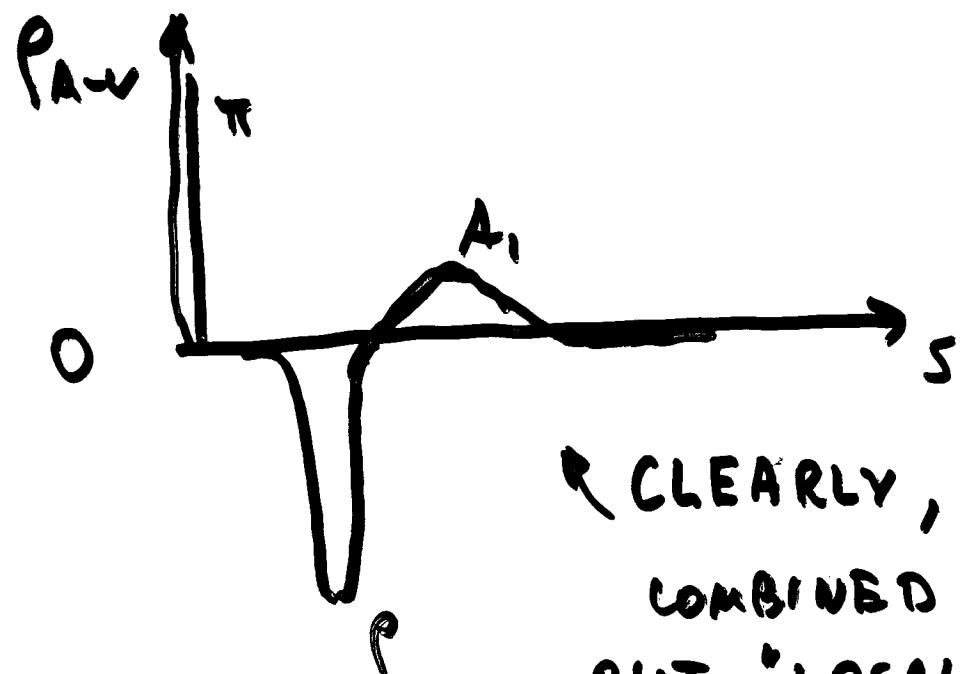
TAKE A-V (OR V-A) COMBINATION
 (e.g., τ -LEPTON DECAY INTO HADRONS)



$$p_A^{\text{quark}} = "1"$$



$$p_V^{\text{quark}} = "1"$$



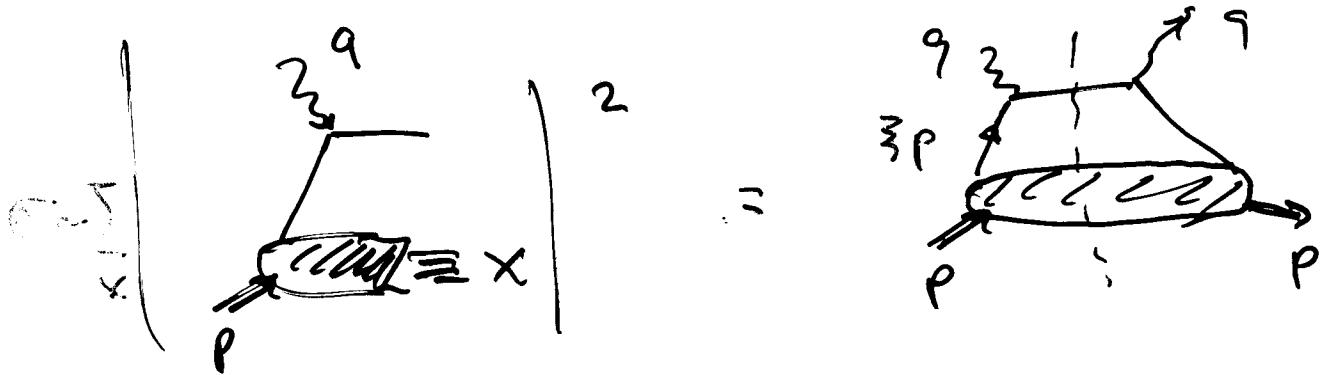
$$p_{A-V}^{\text{quark}} = "0"$$

CLEARLY, THE π, P, A_1
 COMBINED ARE DUAL TO 0
 BUT "LOCAL" DUALITY DOES
 NOT WORK STATE BY STATE

STUDIES OF DUALITY

AT J LAB: DIS IN RESONANCE

REGION.



NACHTMANN VARIABLE

$$(\xi p + q)^2 = 0 \quad \text{- on SHELL QUARK}$$

$$\xi^2 m_p^2 + 2\xi(pq) + q^2 = 0$$

$\hookrightarrow -Q^2$

$$2pq = \frac{Q^2}{x_{Bj'}}$$

$$\xi = \frac{2x_{Bj'}}{1 + \sqrt{1 + 4x_{Bj'}^2 m_p^2 / Q^2}}$$

DGP DUALITY FOR DIS

$$\int_0^1 "z^n" \leftarrow \text{NACHTMANN MOMENTS} F_2(x, Q^2) d\bar{z} = A_n + B_n \frac{1}{Q^2} + C_n \frac{1}{Q^4} + \dots$$

↑ ↑ ↑ ↑
 NACHTMANN/GP TWIST TWIST TWIST
 VARIABLE 2 4 6

$$A_n = \int_0^1 "z^n" F_2^{\text{SCALING}}(\bar{z}, Q^2) d\bar{z}$$

↑ LOG Q^2 -DEPENDENCE
 DUE TO EVOLUTION ONLY

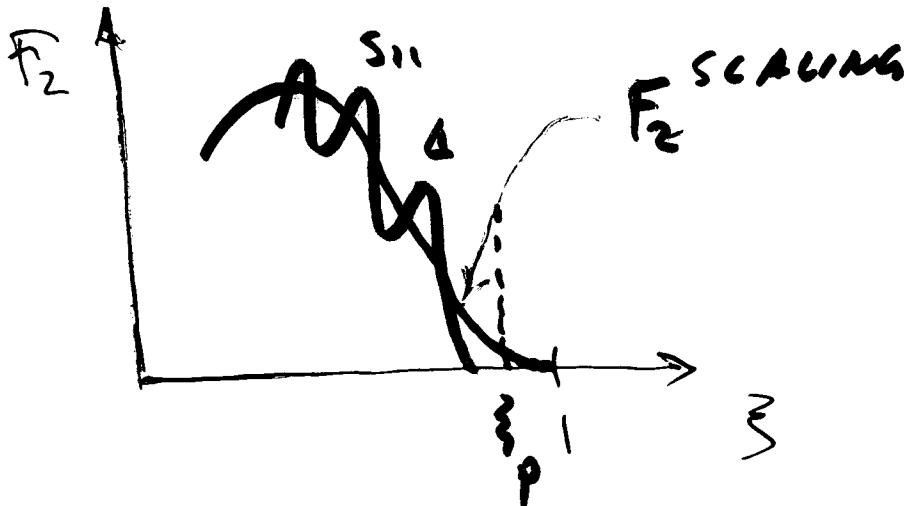
SMALL HIGHER TWISTS FOR $n=0 \rightarrow$

$$\int_0^1 (F_2(x, Q^2) - F_2^{\text{SCALING}}(\bar{z})) d\bar{z} \approx 0.$$

$$\bar{z} = \frac{2x}{1 + \sqrt{1 + 4x^2 m_n^2 / Q^2}}$$

ELASTIC POINT : $x=1 \rightarrow$

$$\bar{z}_p = \frac{2}{1 + \sqrt{1 + 4m_n^2 / Q^2}} < 1$$



NUCLEON ("ELASTIC") CONTRIBUTION

IS DUAL TO SOME INTERVAL

BETWEEN $\bar{z} = 1$ AND \bar{z}_0

Corresponding to

$$w_0^2 \sim \frac{m_N^2 + m_{\text{NEXT}}^2}{2}$$

IF $m_{\text{NEXT}} = m_\Delta$, THEN $w_0^2 \sim 1.2 \text{ GeV}^2$

IF $m_{\text{NEXT}} \sim 1.5 \text{ GeV}$ THEN $w_0^2 \sim 1.6 \text{ GeV}^2$

QCD SUM RULES SUGGEST THAT
DUALITY WORKS FOR CHANNELS WITH
FIXED QUANTUM NUMBERS

$$W_2 \sim \sigma^{1/2} + \sigma^{3/2}$$

$$G_1 \sim \sigma^{1/2} - \sigma^{3/2}$$

$$\begin{array}{c} \downarrow \\ \text{nucleon} \\ + \text{delta} \end{array} \quad \begin{array}{c} \downarrow \\ \text{delta} \end{array}$$

TAKE $(W_2 - G_1)$ - delta is the 1st resonance



" Δ " IS DUAL TO $F_2 - G_1$

SINCE $G_1 \sim F_2$ FOR $x \rightarrow 1$

$F_2 - G_1$ HAS AN EXTRA $\sim(1-x)$ SUPPRESSION

IF DUALITY WORKS FOR Δ IN $F_2 - G_1$,
 $p \rightarrow \Delta$ FORM FACTOR SHOULD FALL
FASTER THAN $p \rightarrow p$ FORM FACTOR
(CALCULATIONS IN PROGRESS)

EXTRACTING $G_n^P(Q^2)$ FROM JLAB DATA IN RESONANCE REGION & DUALITY ASSUMPTION

"DUALITY CURVE"

$$F_2^{\text{DUAL}}(\xi) \approx \xi^{1.12} (1-\xi)^{2.62} (2.7 - 21(1-\xi) + 88(1-\xi)^2 - 131(1-\xi)^3 + 67(1-\xi)^4)$$

(I. NICULESCU)

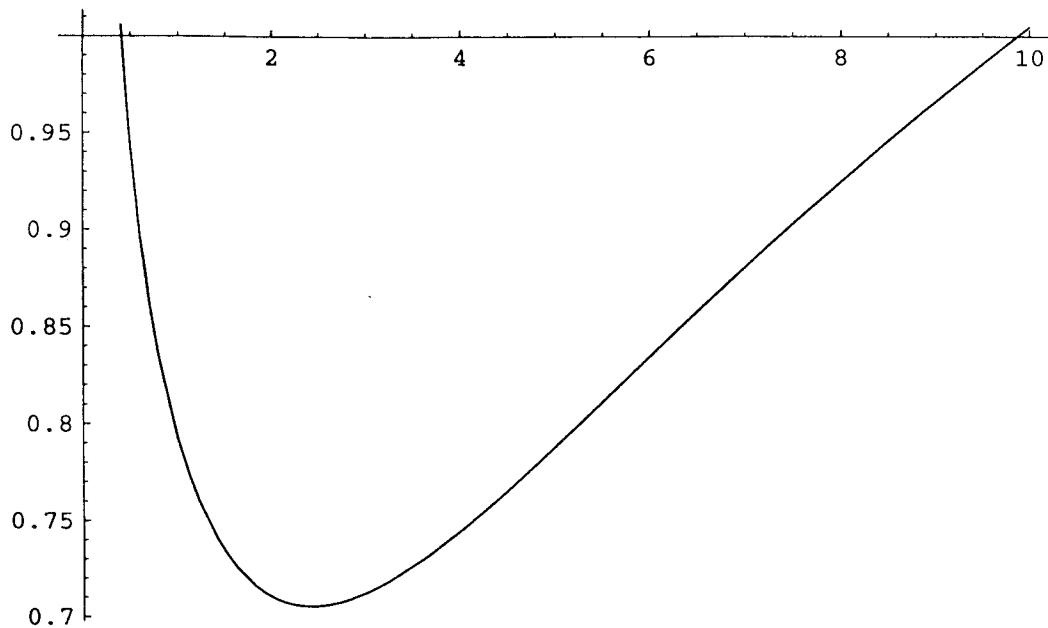
$$F_2^{\text{ELASTIC}}(\xi) = \delta(\xi - \xi_p) \frac{1 + \frac{4m^2}{\pi^2 Q^2}}{1 + \frac{4m^2}{Q^2}} \frac{\xi_p^2}{2 - \xi_p} \left(G_n^P(Q^2)\right)^2$$

$$\int_{\xi_0}^1 F_2^{\text{DUAL}}(\xi) d\xi = \int_{\xi_0}^1 F_2^{\text{ELASTIC}}(\xi)$$

$$\xi_0 = \frac{2x_0}{1 + \sqrt{1 + 4x_0^2 m^2 / Q^2}} , \quad x_0 = \frac{1}{\frac{w^2 - h^2}{Q^2} + 1}$$

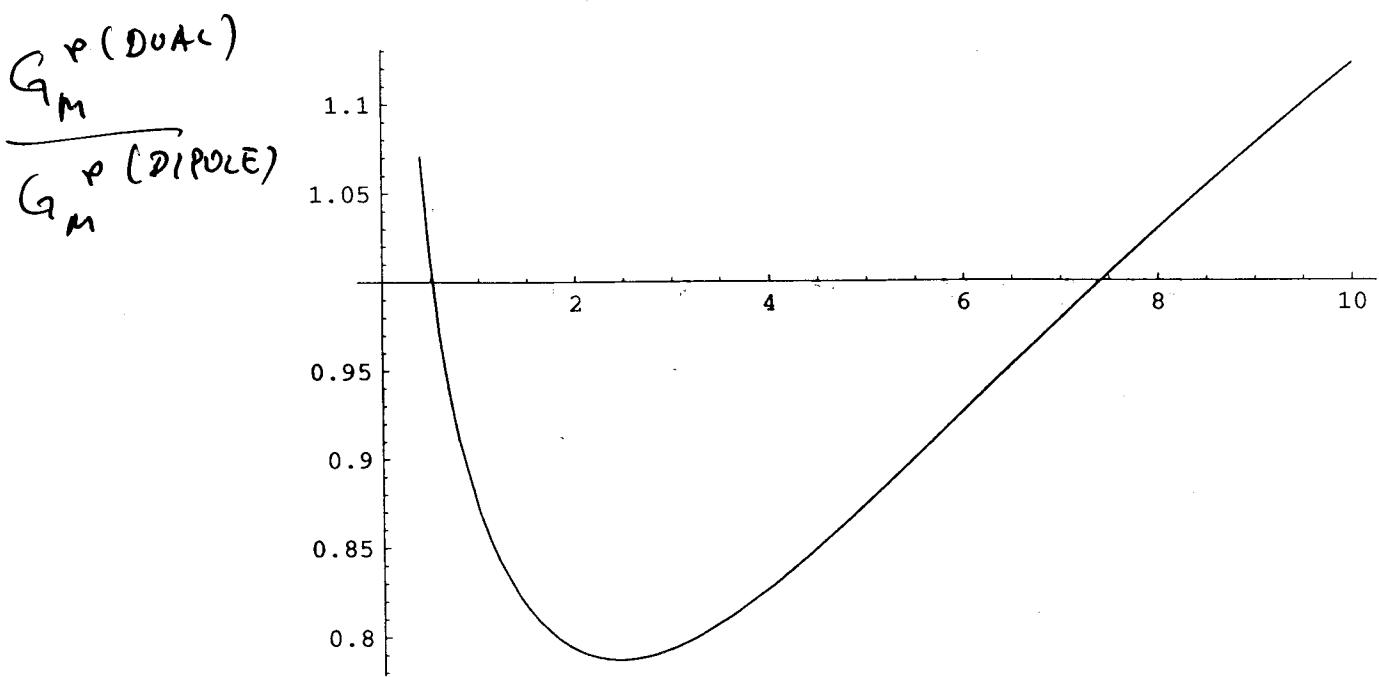
NUCLEON DUALITY INTERVAL

$\frac{G_M^P(\text{DUAL})}{G_M^P(\text{DIPOLE})}$

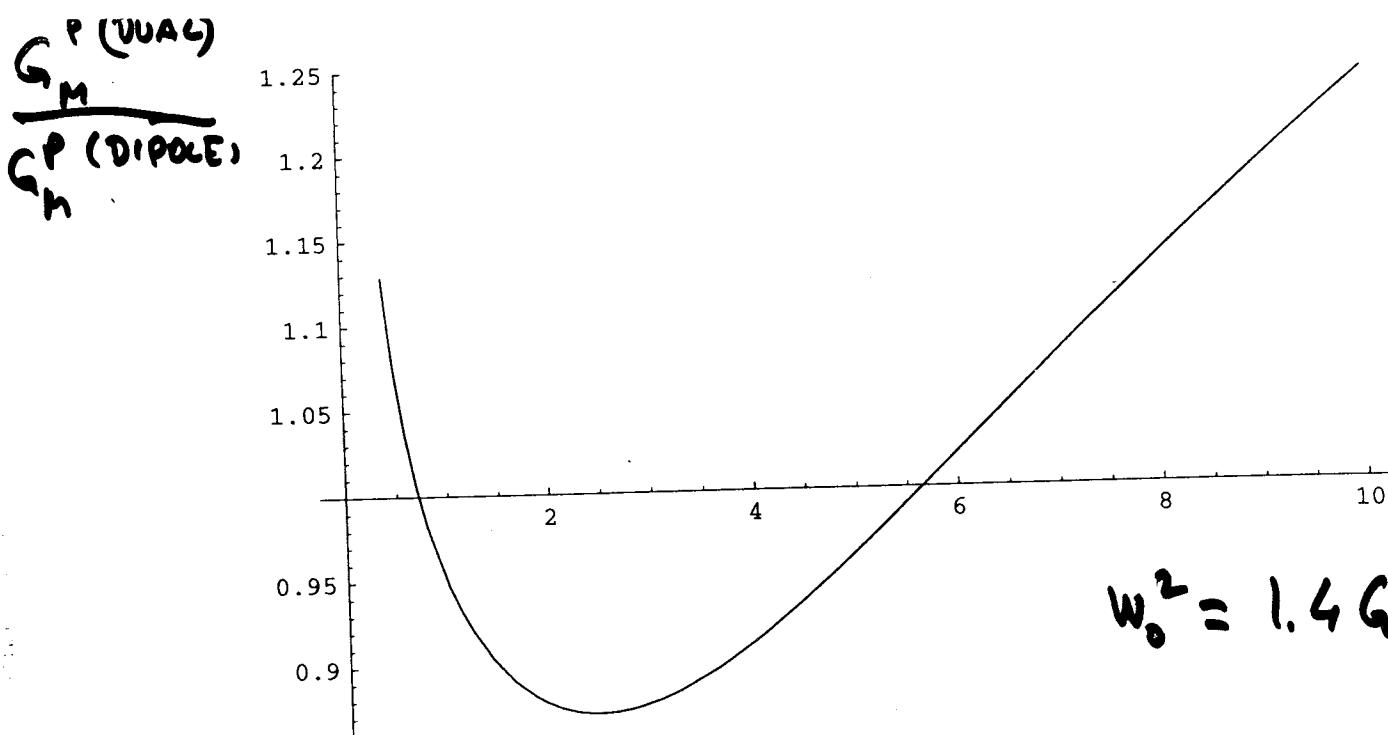


$$W_0^2 = 1.2 \text{ GeV}^2$$

$$\approx (m_\nu + m_\pi)^2$$



$$w^2 = 1.3 \text{ GeV}^2$$

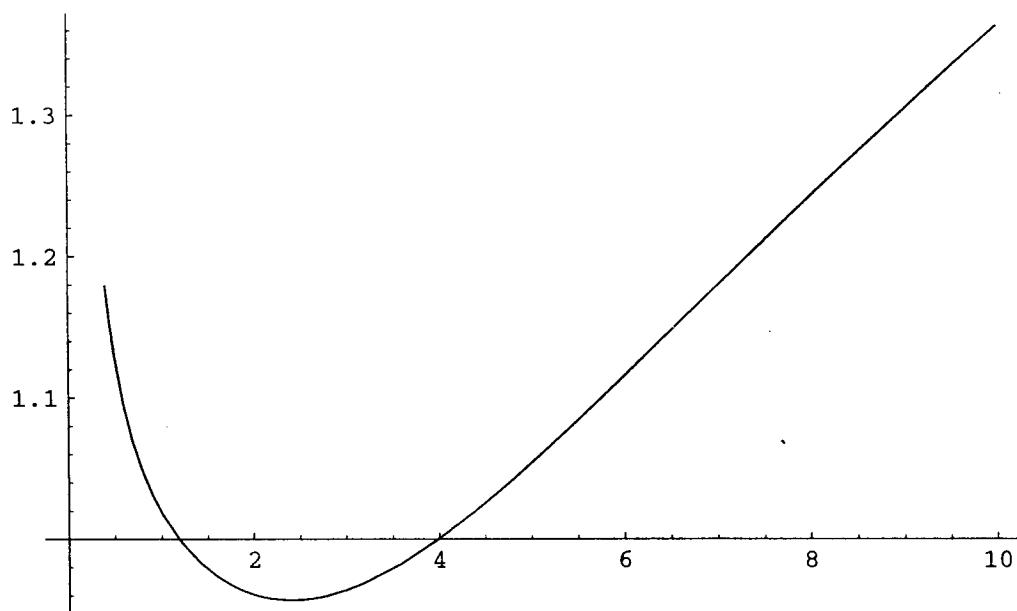


Out[16]= - Graphics -

$$G_M^P(\text{DIPOLE}) = \frac{2.79}{\left(1 + \frac{Q^2}{0.71 \text{ GeV}^2}\right)^2}$$

$$\frac{g_m}{g_m}^{LD}$$

"ex"



$$W^2 \approx 1.5 \text{ GeV}^2$$

CONCLUSIONS

- DIS DATA CAN BE / SHOULD BE USED IN GLOBAL PARAMETERIZATIONS & FOR STUDY OF EVOLUTION AT LOW Q^2 AND LARGE x
- POLARIZED DIS IN THE RESONANCE REGION CAN PROVIDE INTERESTING INFORMATION ON THE CHARACTER OF DUALITY
- SEPARATE $\sigma^{1/2}$ AND $\sigma^{3/2}$ IN FUTURE STUDIES OF LOCAL DUALITY